



**EU - SOUTH KOREA – Joint Researchers Forum
on Semiconductors**



Two-dimensional materials for next generation non-volatile memories

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Brussels (Belgium)

March 25-26, 2024

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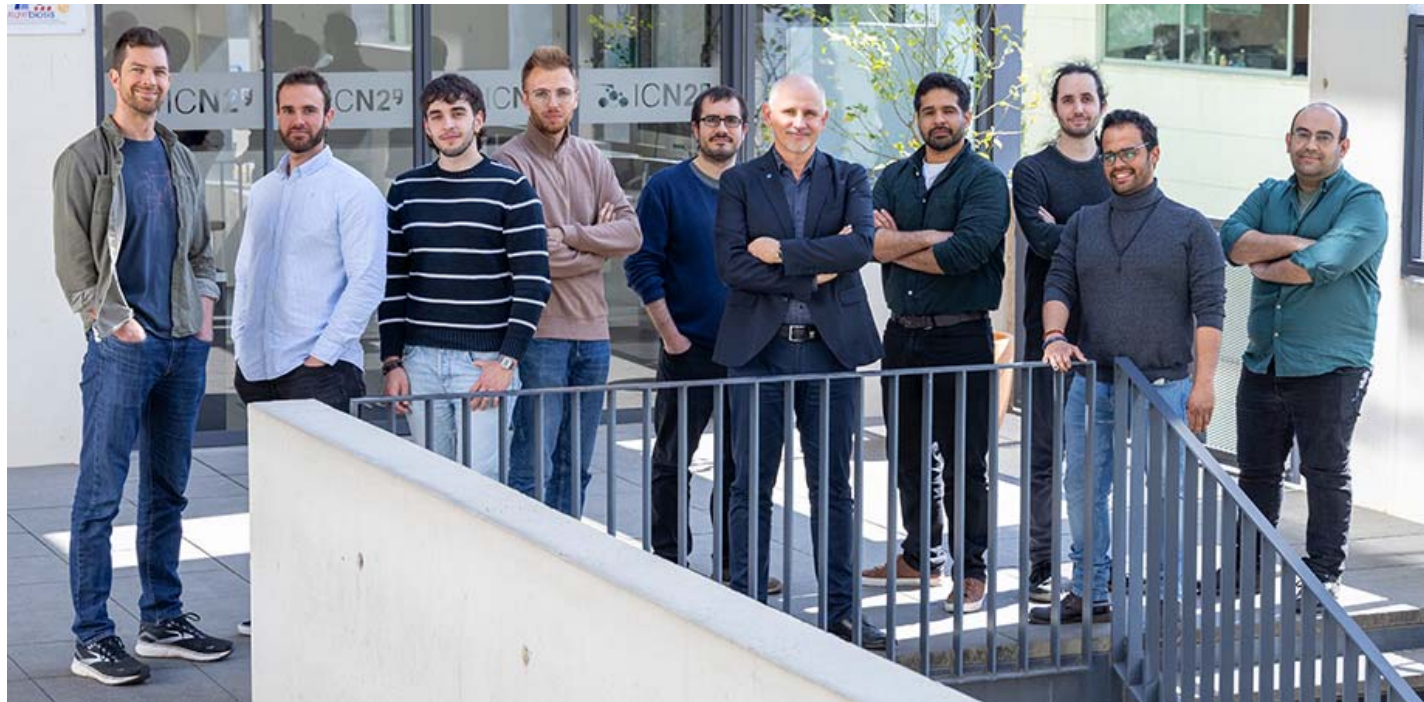
Jose H Garcia



Barcelona, Spain

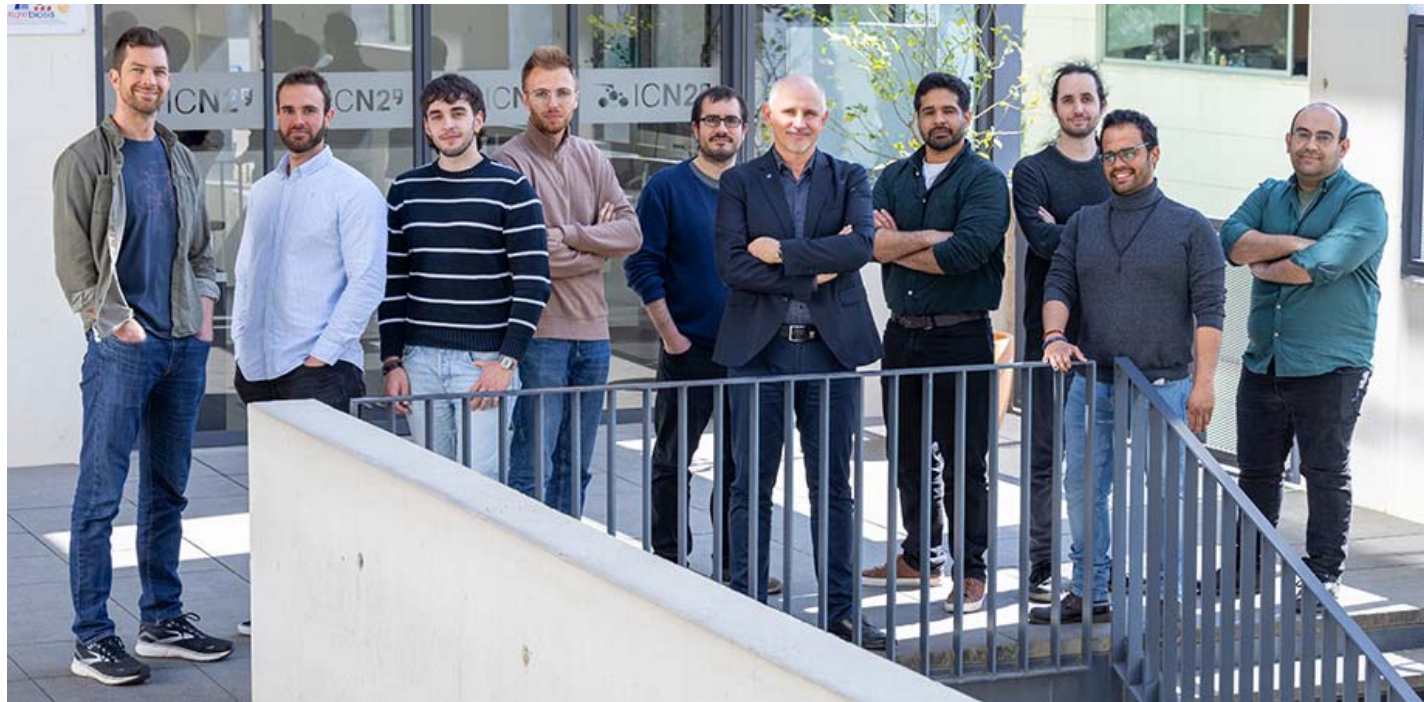
Advanced quantum mechanical simulations

- Disordered materials:
 - Amorphous dielectrics
 - Nonvolatile SOT-MRAM
- Topological quantum devices.
- Quantum device simulations



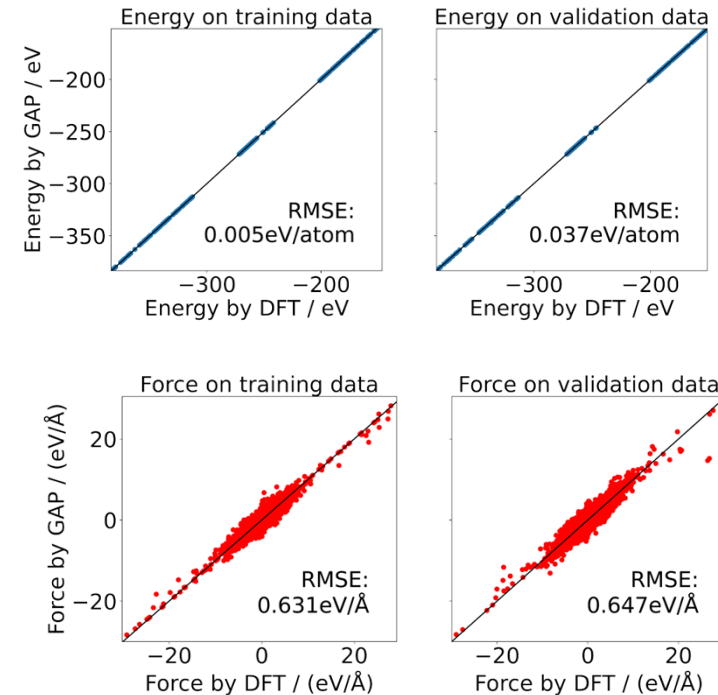
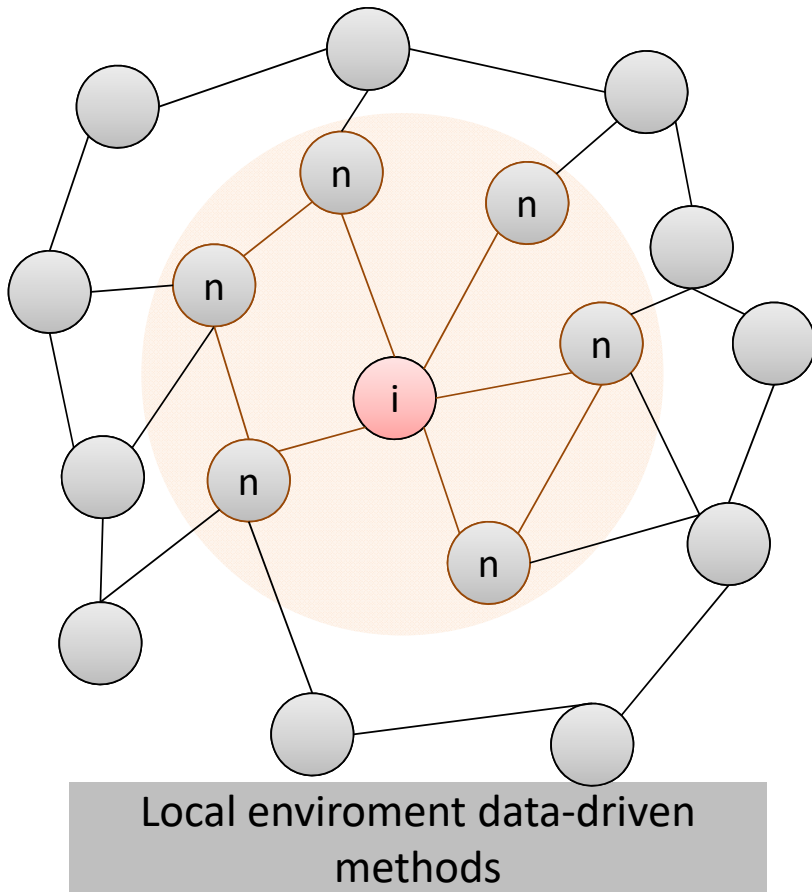
Advanced quantum mechanical simulations

- Disordered materials:
 - Amorphous dielectrics
 - Nonvolatile SOT-MRAM
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Multiscale data-driven simulations

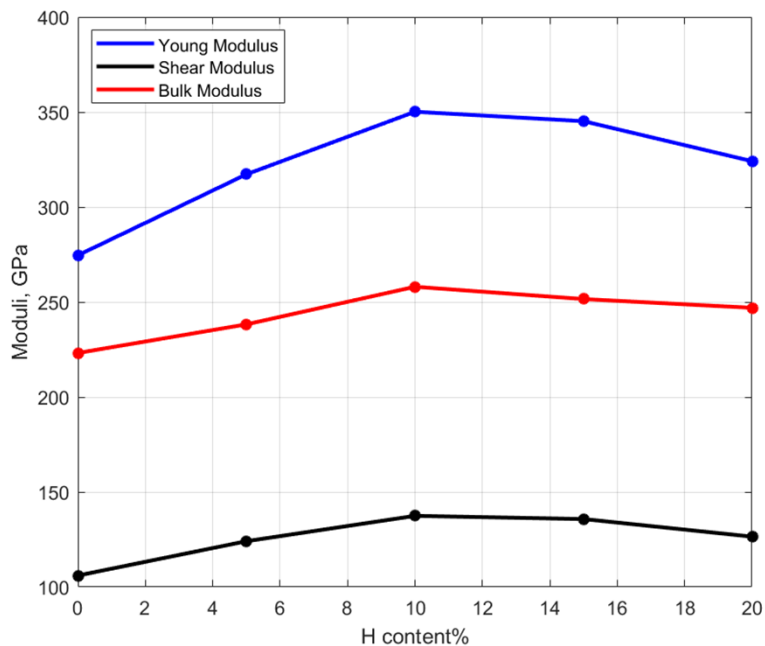
- Machine learning semiclassical potentials.
- Predictive disordered Hamiltonians.
- Linear-Scaling Quantum Transport simulation: Conductivity, Torkance, Optical absorption, etc.



Experiment (hBN): 865 ± 73 Gpa

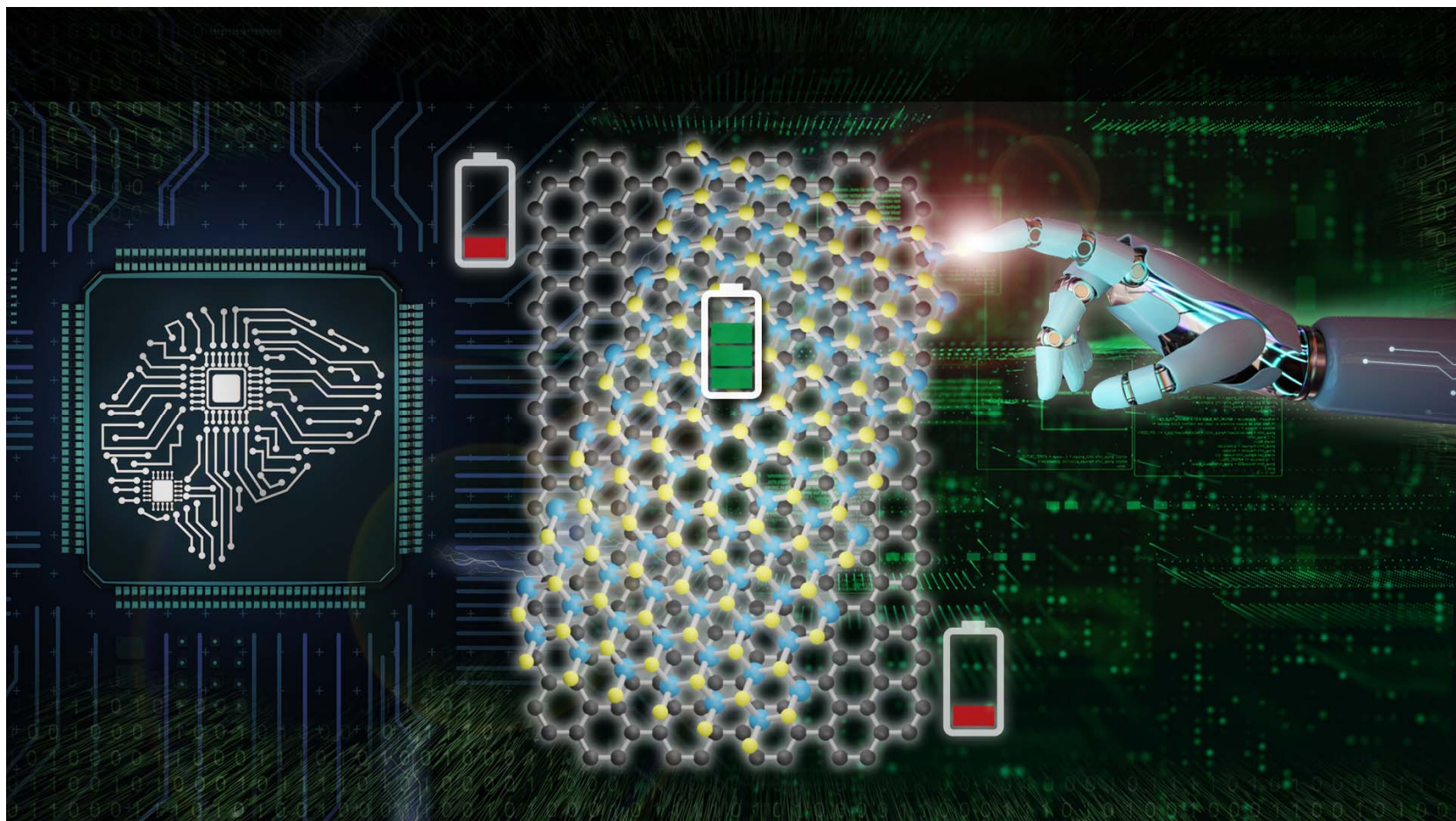
Tersoff potential (hBN): 610 ± 22 GPa

Machine Learning Potential (hBN) 800 ± 15 GPa

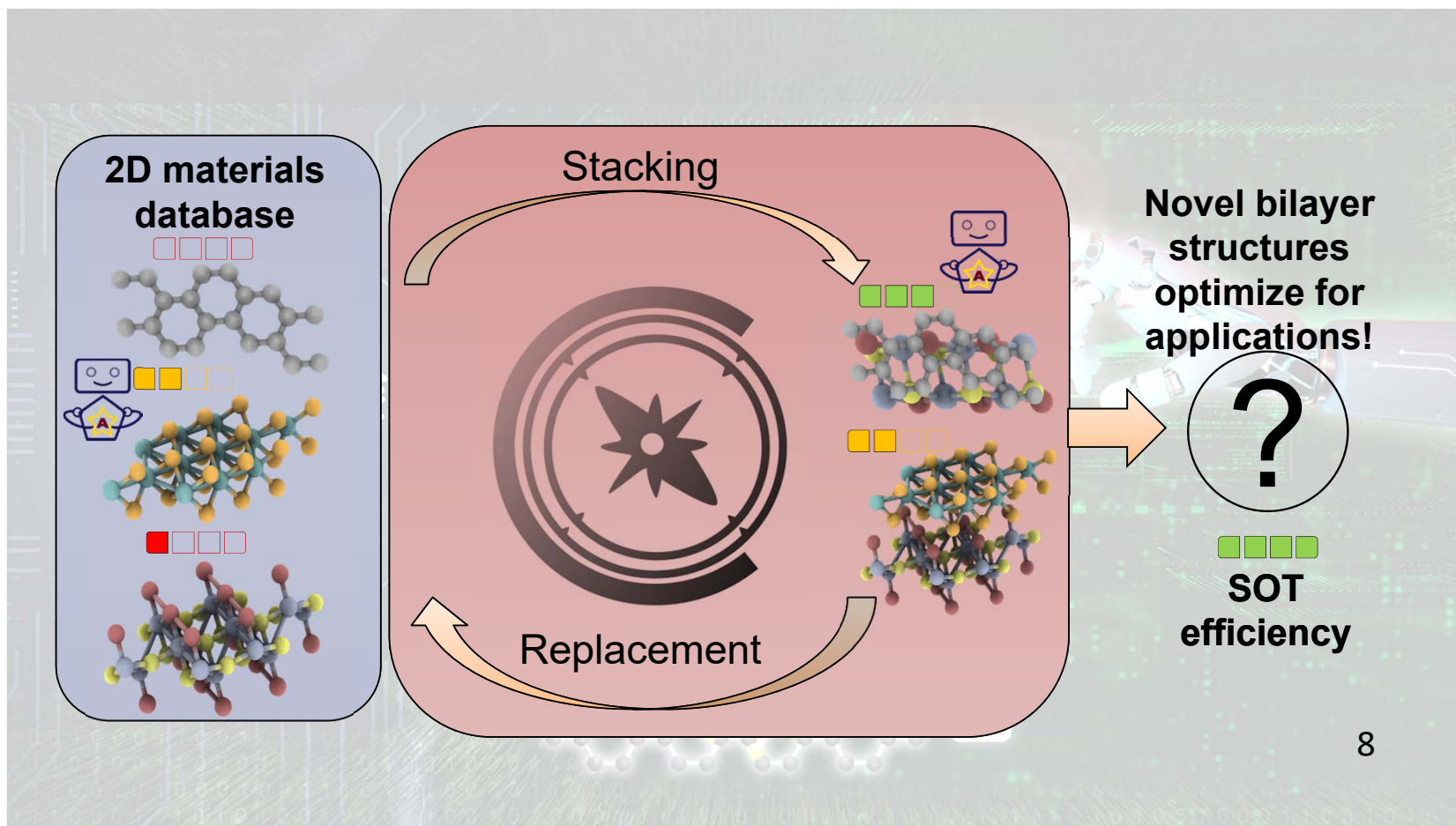


Ultralow-dielectric-constant amorphous boron nitride
Seokmo Hong, et al & Hyeon Suk Shin Nature **582**,
511–514 (2020)

Combine
heuristic
optimization and
AI for spin-orbit
memories



Combine heuristic optimization and AI for spin-orbit memories



$$\chi_{\alpha,\beta}^A(\mu, T) = \int_{-\infty}^{\infty} d\varepsilon f(\mu, \varepsilon, T) \text{Tr} \left[\frac{dG^+(H, \varepsilon)}{d\varepsilon} A_{\alpha} \text{Im}[G^+(H, \varepsilon)] j_{\beta} - \text{Im}[G^-(H, \varepsilon)] A_{\alpha} \frac{dG^-(H, \varepsilon)}{d\varepsilon} j_{\beta} \right]$$

- A_{α}, j_{β} operator matrices
- H Hamiltonian matrix
- $G^{\pm}(H, \varepsilon)$ Green's functions

$N = 10^6 - 10^9$
atoms

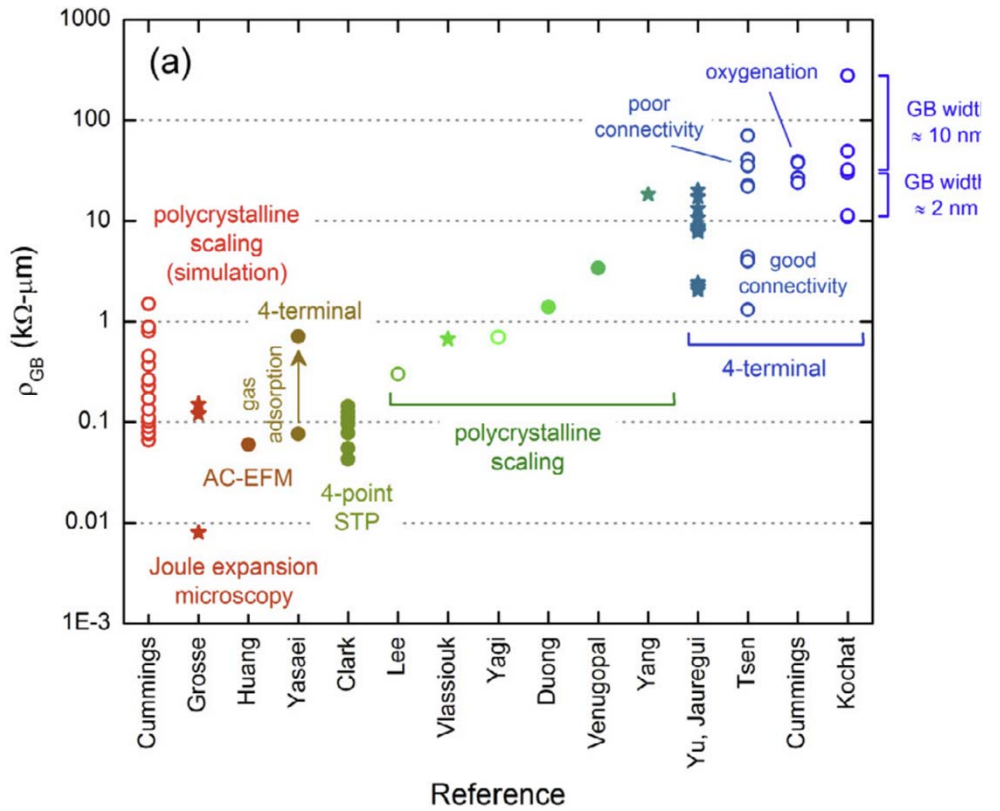
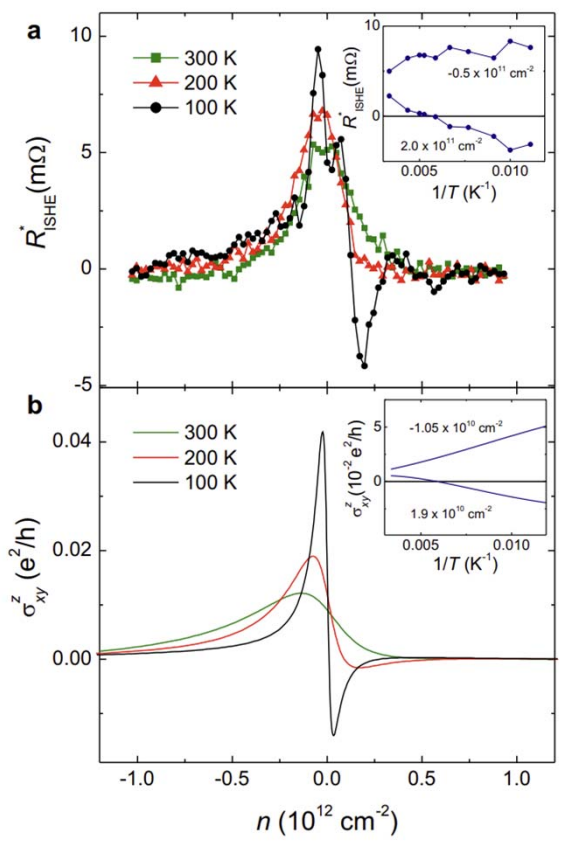
$O(N)$ real-space algorithms
 $O(N^3) \longrightarrow O(N)$
 1 billion years \longrightarrow 1 month



Z. Fan, Jose H.Garcia, A W. Cummings et al, Physics Reports 903, 7, 1-69 (2021)
 Jose H. Garcia, Lucian Covaci, Tatiana G. Rappoport , *Phys. Rev. Lett.* 114, 116602 (2015)

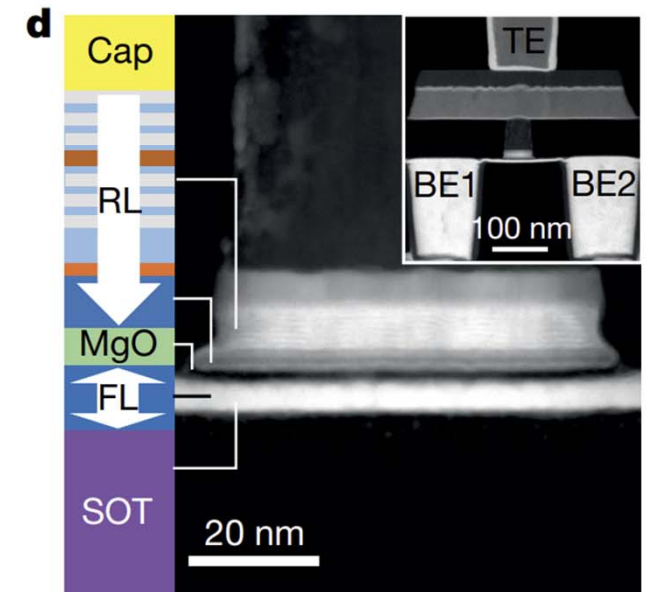
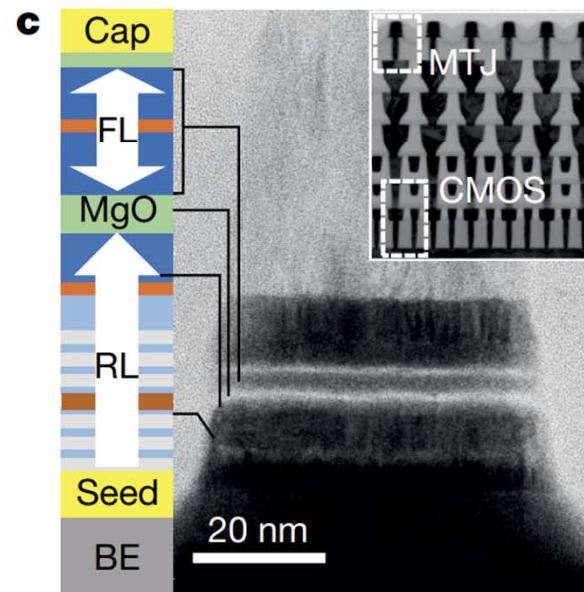
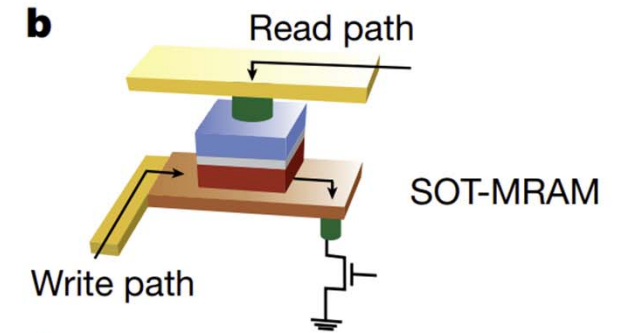
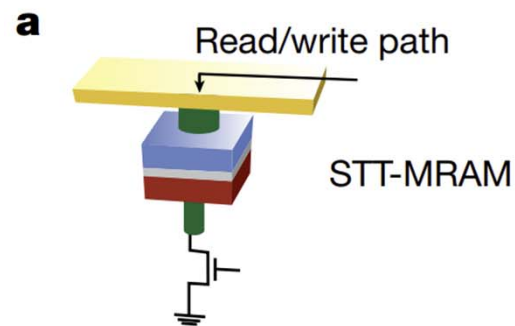
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 Jose H Garcia , ICN2

Linear-scaling quantum transport methodologies



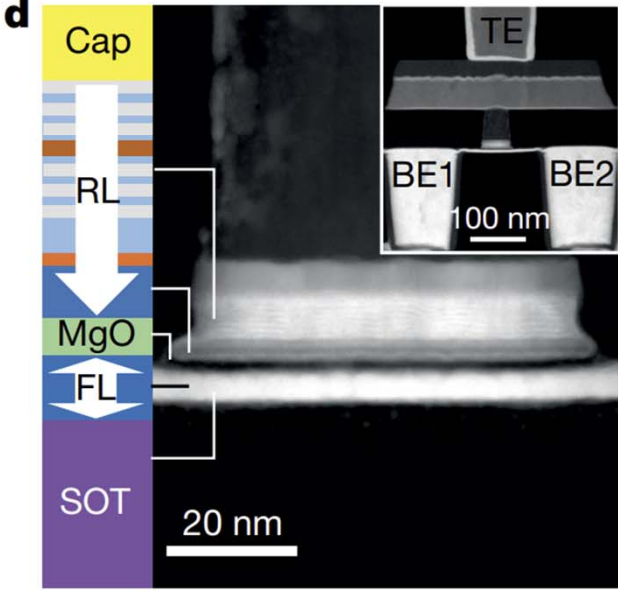
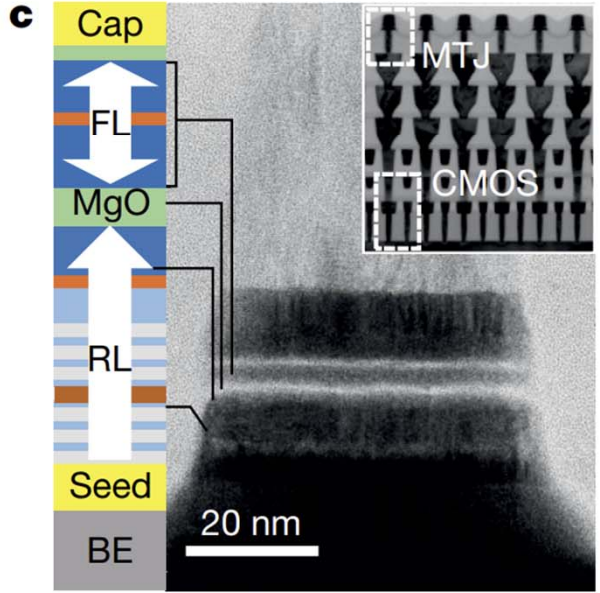
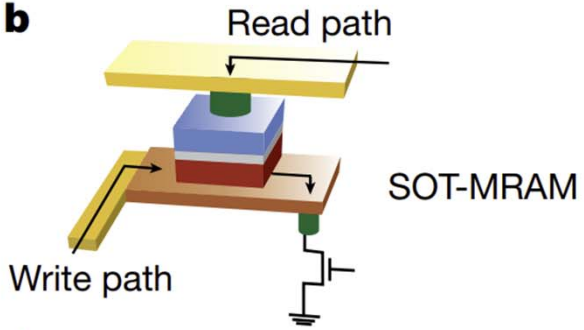
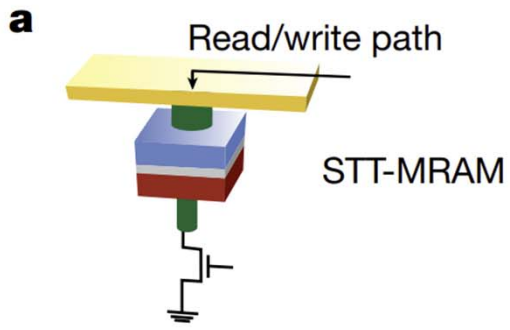
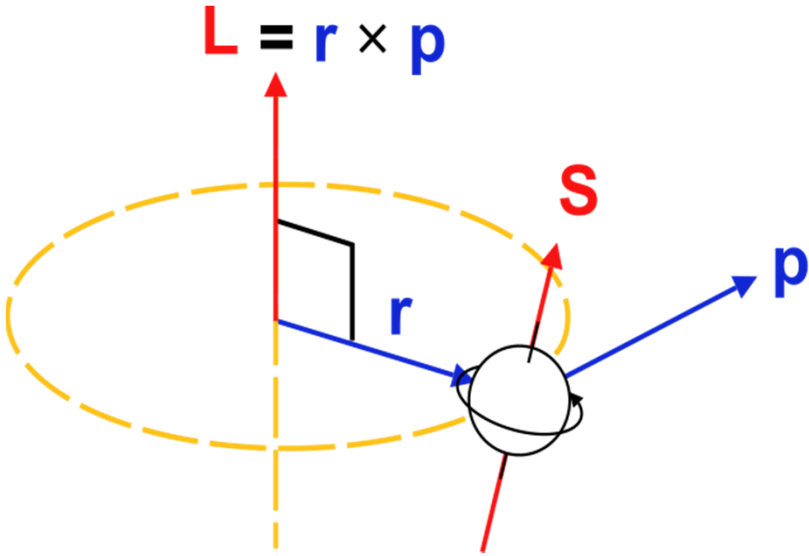
Magnetic Random Access Memories (MRAMs)

Why to use spin-orbit coupling?

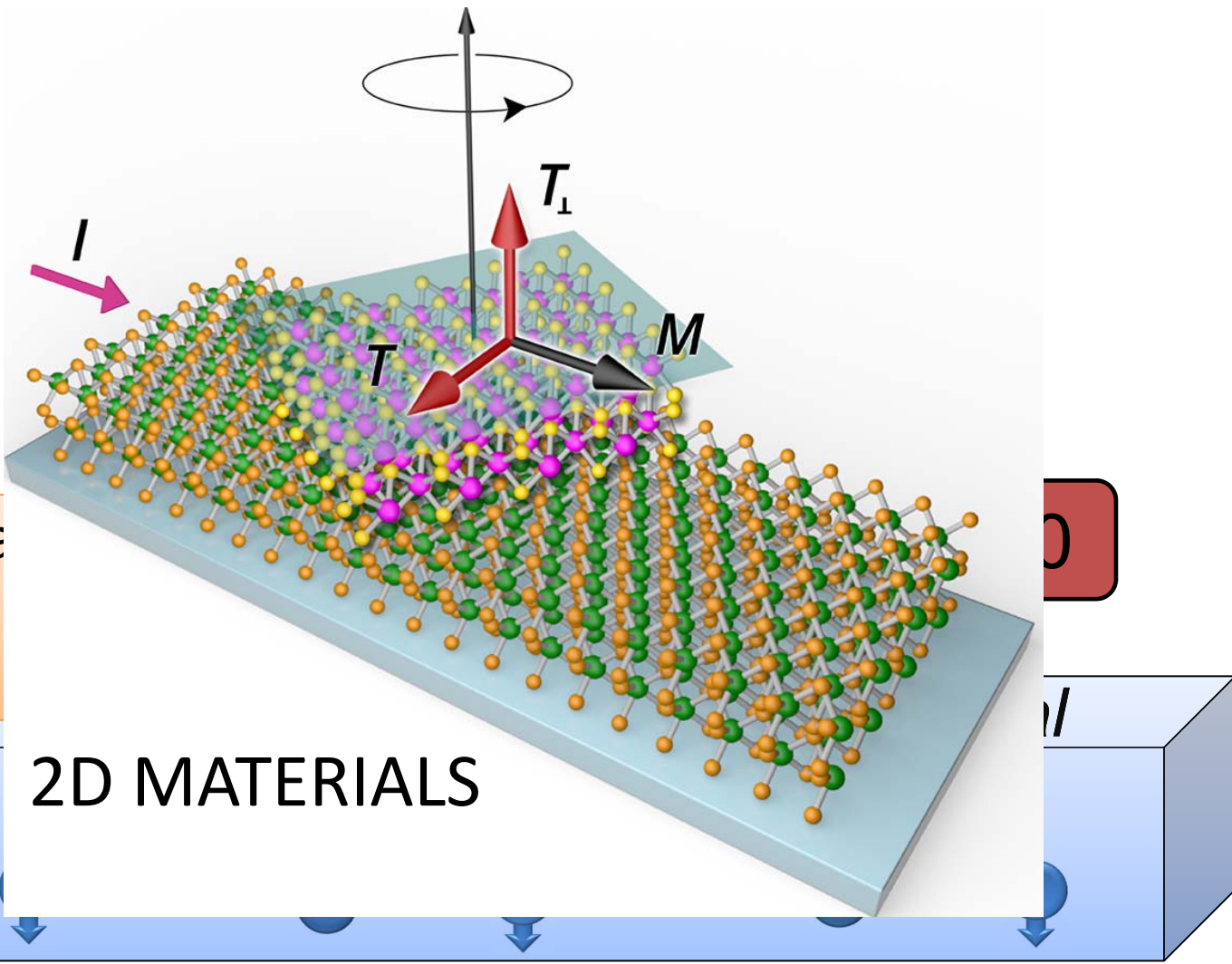


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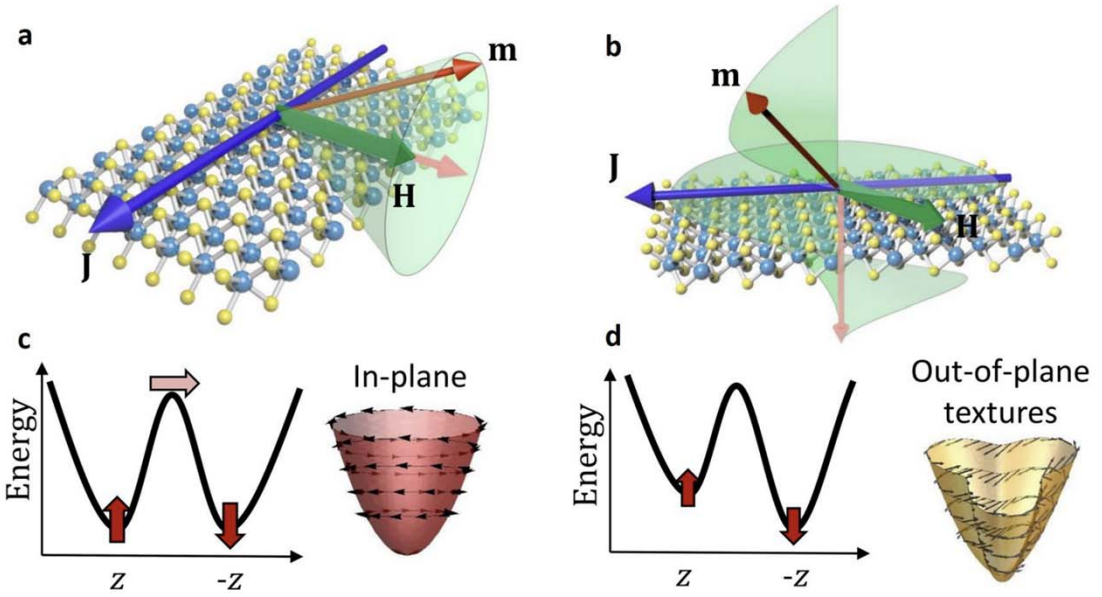
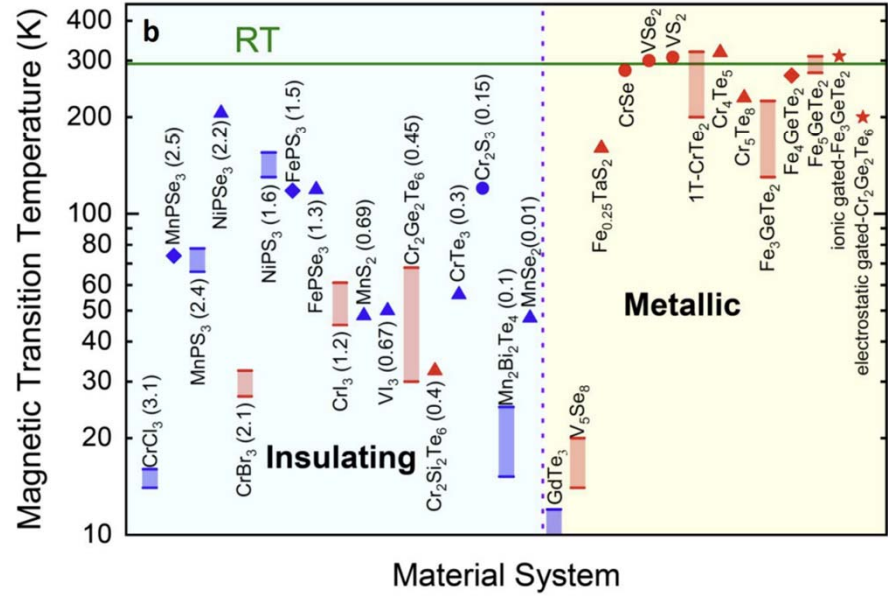
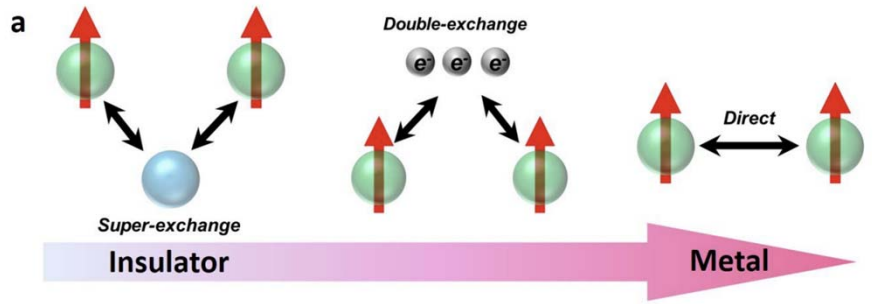


HOW DOES SOT MRAM WORK?



Interfacial effect that constrains to ultrathin systems

Two-dimensional magnets



H Kurebayashi, Jose H. Garcia, Safe Khan, Jairo Sinova & Stephan Roche
 Nature Reviews Physics volume 4, pages150–166 (2022)

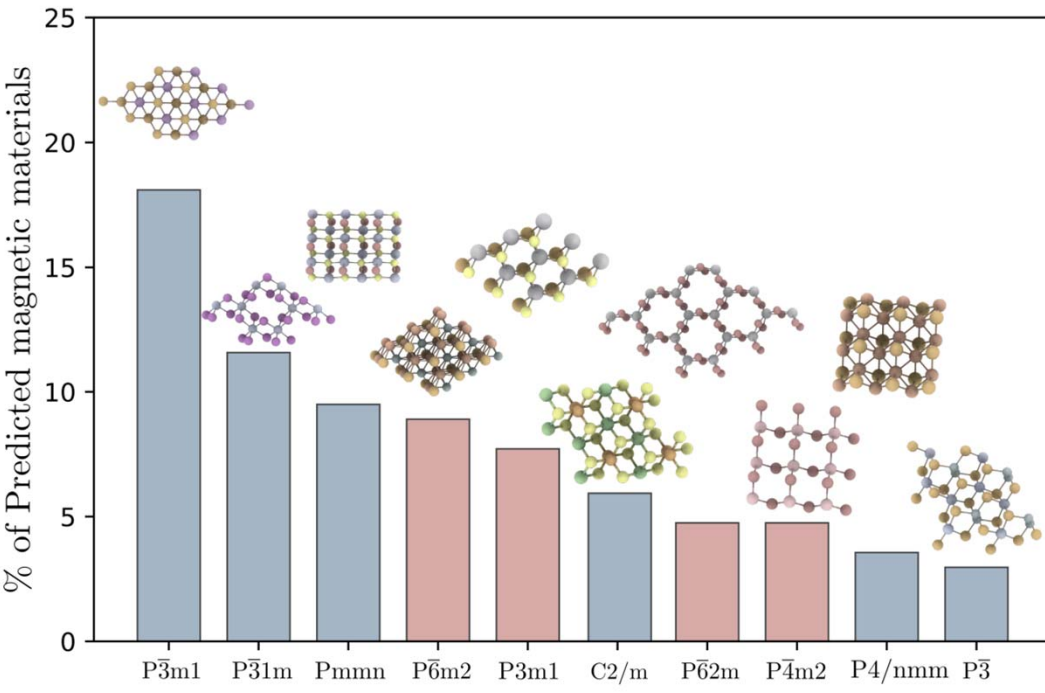
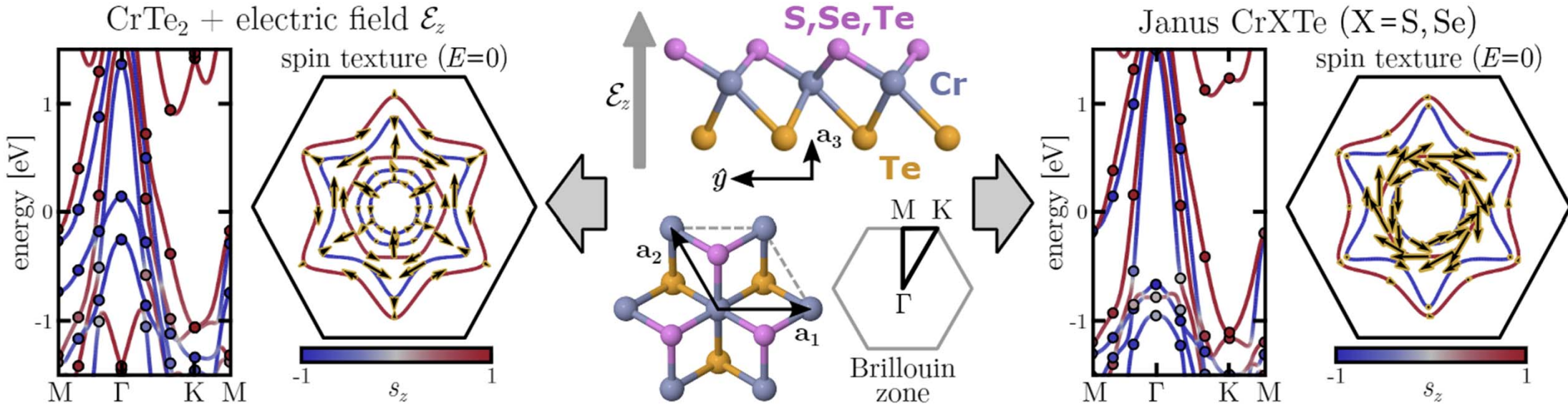
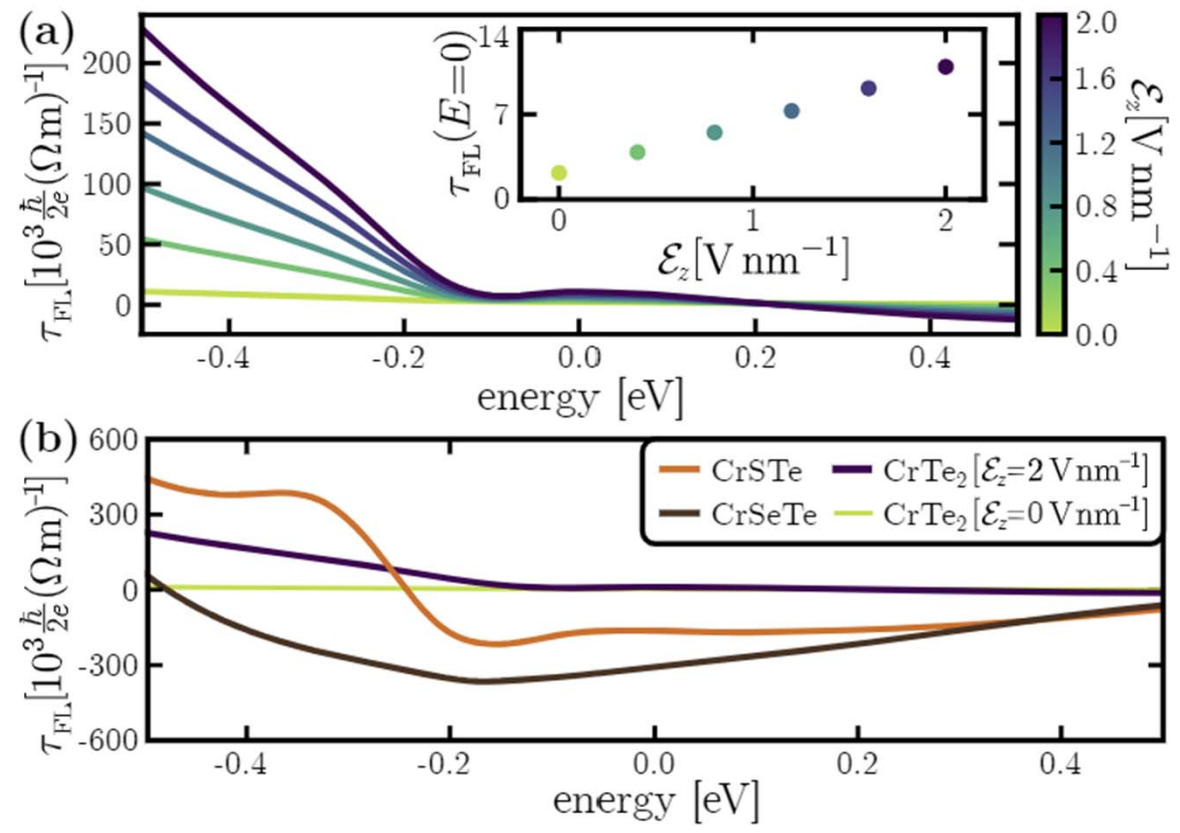


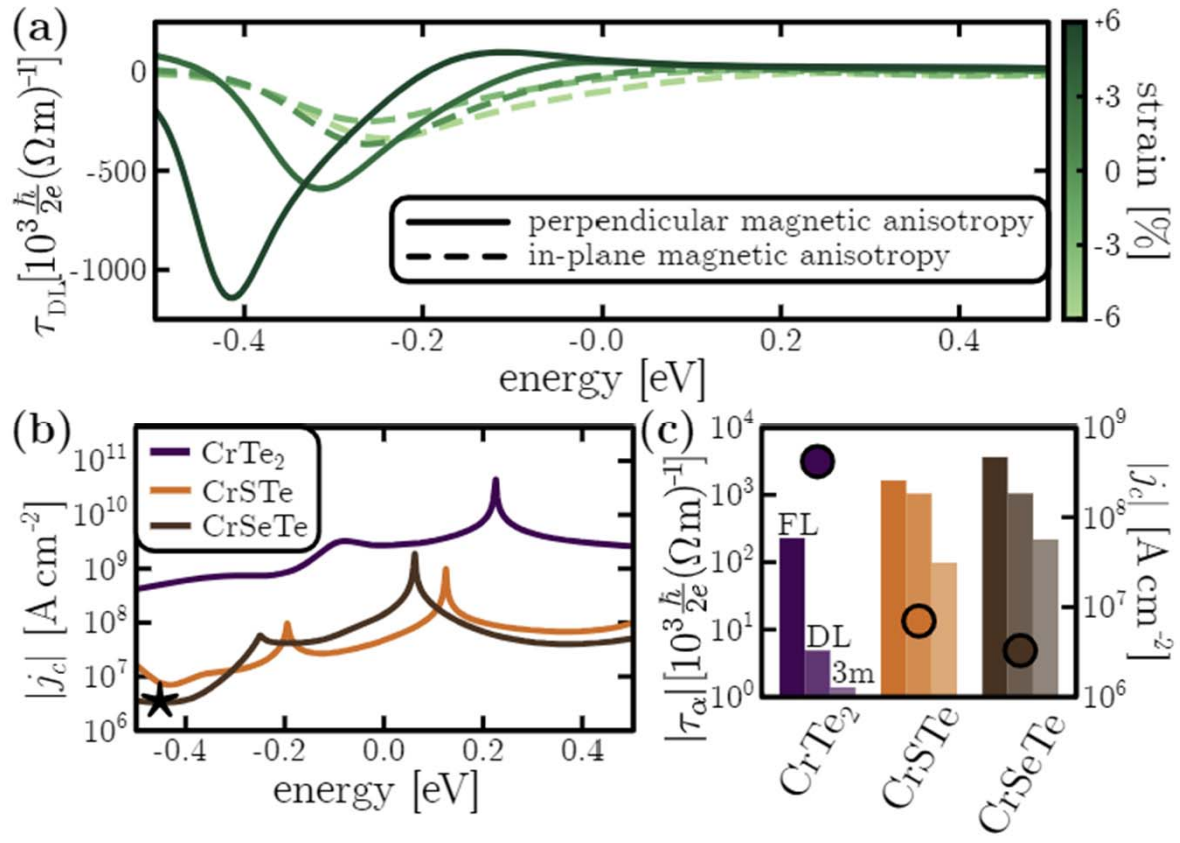
Table 1

Repr. Mat.	Space Group	Crystal System	H^e	H^o
MX_3^*	P62m (189)	Hexagonal	$\tau_{FL} \mathbf{J} \times \hat{\mathbf{z}}$	$\tau_{DL} \mathbf{m} \times (\mathbf{J} \times \hat{\mathbf{z}}) + \tau_z m_z \mathbf{J} + \tau_{ani} \nabla_m [J_y m_x m_y + J_x (m_y^2 - m_x^2)/2]$
$Fe_3GeTe_2^*$	P6m2 (187)			$\tau_{DL} \mathbf{m} \times (\mathbf{J} \times \hat{\mathbf{z}}) + \tau_z m_z \mathbf{J} + \tau_{ani} \nabla_m [J_x m_x m_y + J_y (m_x^2 - m_y^2)/2]$
$MnBi_2Te_2$	P3m1 (164)	Trigonal	$\tau_{FL} \mathbf{J} \times \hat{\mathbf{z}}$	$\tau_{DL} \mathbf{m} \times (\mathbf{J} \times \hat{\mathbf{z}}) + \tau_z m_z \mathbf{J} + \tau_{ani} \nabla_m [J_x m_x m_y + J_y (m_x^2 - m_y^2)/2]$
MX_2^\dagger	P3m1 (156)			
Fe_3GeTe_2	R3m (147)			
CrI_3	P31m (162)			$\tau_{DL} \mathbf{m} \times (\mathbf{J} \times \hat{\mathbf{z}}) + \tau_z m_z \mathbf{J} + \tau_{ani} \nabla_m [J_y m_x m_y + J_x (m_y^2 - m_x^2)/2]$
$CrGeTe_3$	P3 (147)	Trigonal	$\tau_{FL}^I \mathbf{J} \times \hat{\mathbf{z}} + \tau_{FL}^{II} \mathbf{J}$	$\tau_{DL}^I \mathbf{m} \times (\mathbf{J} \times \hat{\mathbf{z}}) + \tau_z m_z \mathbf{J} + \tau_{ani}^I \nabla_m \left[J_y m_x m_y + \frac{J_x (m_x^2 - m_y^2)}{2} \right] + \tau_{ani}^{II} \nabla_m \left[J_y m_x m_y + \frac{J_x (m_x^2 - m_y^2)}{2} \right] + \tau_o^{II} [(\mathbf{J} \times \mathbf{m}) \cdot \hat{\mathbf{z}}] \hat{\mathbf{z}}$
$FeTe^\dagger$	P4/nmm (129)	Tetragonal	$\tau_{FL} \mathbf{J} \times \hat{\mathbf{z}}$	$\tau_o \mathbf{m} \times (\mathbf{J} \times \hat{\mathbf{z}}) + \tau_z m_z \mathbf{J}$
CrI_2	P4m2 (115)	Orthorhombic	$\begin{pmatrix} 0 & \tau_{FL}^{xy} \\ \tau_{FL}^{yx} & 0 \end{pmatrix} \mathbf{J}$	$(\mathbf{M} \cdot \hat{\mathbf{z}}) \begin{pmatrix} \tau_z^x & 0 \\ 0 & \tau_z^y \end{pmatrix} \mathbf{J} + \mathbf{M} \cdot \begin{pmatrix} \tau_o^x & 0 \\ 0 & \tau_o^y \end{pmatrix} \mathbf{J} \hat{\mathbf{z}}$
$CrSBr$	Pnmm (59)	Orthorhombic	$\begin{pmatrix} 0 & \tau_{FL}^{xy} \\ \tau_{FL}^{yx} & 0 \end{pmatrix} \mathbf{J}$	$(\mathbf{M} \cdot \hat{\mathbf{z}}) \begin{pmatrix} \tau_z^x & 0 \\ 0 & \tau_z^y \end{pmatrix} \mathbf{J} + \mathbf{M} \cdot \begin{pmatrix} \tau_o^x & 0 \\ 0 & \tau_o^y \end{pmatrix} \mathbf{J} \hat{\mathbf{z}}$
$NiPS_3$	C2/m (12)	Monoclinic	$\begin{pmatrix} 0 & \tau_{FL}^{xy} & 0 \\ \tau_{FL}^{yx} & 0 & 0 \\ 0 & \tau_{FL}^{yz} & 0 \end{pmatrix} \mathbf{J}$	$\begin{pmatrix} M_x \tau_o^{xxx} + M_z \tau_o^{zxx} & M_y \tau_o^{yyx} & 0 \\ M_y \tau_o^{yxy} & M_x \tau_o^{xxy} + M_z \tau_o^{zyy} & 0 \\ M_z \tau_o^{zxx} & 0 & 0 \end{pmatrix} \mathbf{J} + \mathbf{M} \cdot \begin{pmatrix} \tau_o^{xx} & 0 & 0 \\ 0 & \tau_o^{yy} & 0 \\ 0 & 0 & 0 \end{pmatrix} \mathbf{J} \hat{\mathbf{z}}$

H Kurebayashi, Jose H. Garcia, Safe Khan, Jairo Sinova & Stephan Roche
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- Two-dimensional materials are showing great promise for highly optimized non-volatile magnetic random access memories (MRAM), thanks to their interfacial characteristics, high tunability, and compatibility with existing technologies. Currently, Cr-based magnetic Janus materials appear to be among the most promising candidates, offering critical switching currents on par with those of the leading metal/magnetic interfaces.



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This work was supported through the I+D+i project PID2022-138283NB-I00, financed by MICIU/AEI/10.13039/501100011033/ and "FEDER Una manera de hacer Europa".

THANK YOU



This project has received funding from the European Union's Horizon Europe research and innovation programme under GA N° 101092562

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